

## SOME CONFORMAL TRANSFORMS

Write  $t = u + jv$  and  $z = x + jy$ , where  $u, v, x$  and  $y$  are real. Let  $A$  be the region in the  $t$ -plane where  $v > 0$ , and let  $B$  be the region in the  $z$ -plane where  $-1 < x < 1$  and  $y > 0$ .

I claim that there is an invertible conformal mapping from  $A$  to  $B$  given by  $z = 2 \arcsin(t)/\pi$ , or  $t = \sin(\pi z/2)$ . This can either be seen directly, or using the general framework of Schwartz-Christoffel transforms.

For the direct argument, consider the map  $f: B \rightarrow \mathbb{C}$  given by

$$t = f(z) = \sin(\pi z/2) = \sin(\pi x/2) \cosh(\pi y/2) + \cos(\pi x/2) \sinh(\pi y/2)j.$$

When  $z \in B$  we have  $|x| < 1$  (so  $\cos(\pi x/2) > 0$ ), and also  $y > 0$  (so  $\sinh(y) > 0$ ). We conclude that  $v = \text{Im}(t) = \text{Im}(f(z)) > 0$ , so the point  $t = f(z)$  lies in  $A$ , so we actually have a function  $f: B \rightarrow A$ . The bottom edge of  $B$  is the real interval  $[-1, 1]$ , or equivalently, the interval where  $y = 0$  and  $-1 \leq x \leq 1$ . By inspecting the graph of  $f$ , we see that  $f$  carries this interval to itself in an invertible manner. Next, we have

$$\begin{aligned} f(-1 + jy) &= -\cosh(\pi y/2) \\ f(+1 + jy) &= +\cosh(\pi y/2). \end{aligned}$$

The left hand edge of  $B$  is the line where  $x = -1$  and  $y > 0$ . By inspecting the graph of  $\cosh$ , we see that  $\cosh(\pi y/2)$  runs from 1 to  $\infty$  along this line, so  $f$  carries the left edge of  $B$  to the interval  $(-\infty, -1]$  in an invertible manner. Similarly, it carries the right hand edge to  $[1, \infty)$ . Putting all this together, we see that  $f$  carries the boundary of  $B$  invertibly to the real axis, which is the boundary of  $A$ . The general theory of conformal mappings now allows us to conclude that  $f$  carries  $B$  invertibly to  $A$ .

Alternatively, Schwartz-Christoffel theory says that there is an invertible conformal mapping between  $A$  and  $B$ , such that

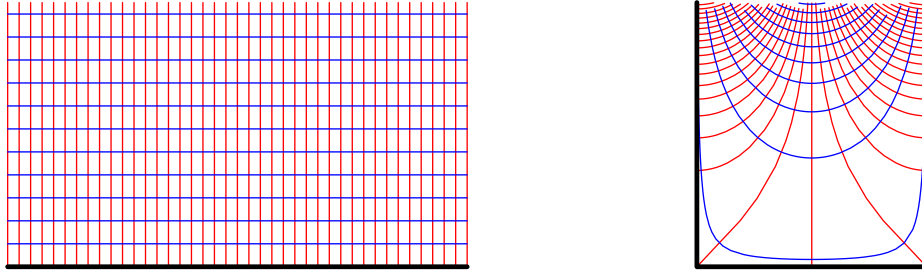
$$\frac{dz}{dt} = K(t - t_0)^{-1/2}(t - t_1)^{-1/2},$$

for some constants  $K, t_0$  and  $t_1$  with  $t_0$  and  $t_1$  real. We are free to make an additional transformation of the form  $t \mapsto \alpha t + \beta$  (with  $\alpha$  and  $\beta$  real and  $\alpha \neq 0$ ). Using this, we may assume that  $t_0 = -1$  and  $t_1 = +1$ , and thus that  $(t - t_0)(t - t_1) = (t + 1)(t - 1) = t^2 - 1$ . This gives

$$\begin{aligned} z &= K \int \frac{dt}{\sqrt{t^2 - 1}} \\ &= Kj \int \frac{dt}{\sqrt{1 - t^2}} \\ &= Kj \arcsin(t) + c. \end{aligned}$$

When  $t = 1$  this gives  $z = Kj\pi/2$ , but the point  $t = t_1 = 1$  must correspond to the corner of  $B$  at  $z = 1$ , so we have  $1 = Kj\pi/2 + c$ . By considering  $t = -1$  in the

same way, we see that  $-1 = -Kj\pi/2 + c$ . Adding these two equations gives  $c = 0$ , and it then follows that  $Kj = 2/\pi$ , so  $z = 2 \arcsin(t)/\pi$  as claimed.

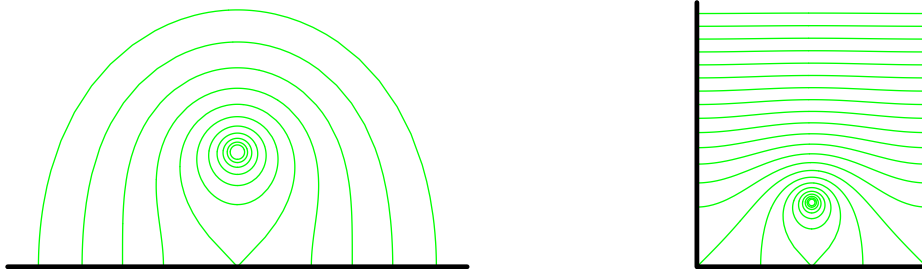


The picture on the left shows the  $t$ -plane, with lines of constant  $u$  marked in red, and lines of constant  $v$  in blue. The picture on the right shows the corresponding curves in the  $z$ -plane.

Now consider a current flowing in a wire passing perpendicularly through our diagram at the point  $t = j$ , and suppose that the boundaries of the regions  $A$  and  $B$  are made of infinitely permeable iron. The method of images says that the resulting magnetic field in  $A$  is the same as the field coming from two equal currents at  $j$  and  $-j$ , with no iron boundary. The complex potential for this field is  $\omega = \log((t-j)(t+j)) = \log(t^2 + 1)$  (or a positive multiple of this). The magnetic flux function is thus

$$\phi = \text{Re}(\omega) = \text{Re}(\log(t^2 + 1)) = \log |t^2 + 1|.$$

The field lines are the lines  $\phi = c$  for constant  $c$ . The equation  $\phi = c$  is equivalent to  $|t^2 + 1| = e^c$ , or  $t^2 + 1 = e^{c+i\theta}$  for some  $\theta$ , so  $t = (1 - e^{c+i\theta})^{1/2}$ . This makes it easy to plot the field lines in the  $t$ -plane, as shown on the left below.



As the two-dimensional magnetic field equations are conformally invariant, the field lines in the  $z$ -plane are obtained by just applying the rule  $z = 2 \arcsin(t)/\pi$  to the field lines in the  $t$ -plane. The result is shown on the right in the diagram. Note that the current is located at the point

$$z = 2 \arcsin(j)/\pi = (2 \operatorname{arcsinh}(1)/\pi)j \simeq 0.56j.$$