

Divergence and curl

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Div and curl of matrix fields

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$$\operatorname{div}(\mathbf{u}) = \frac{\partial}{\partial x}(ax + by + cz) + \frac{\partial}{\partial y}(dx + ey + fz) + \frac{\partial}{\partial z}(gx + hy + iz)$$

$$= a + e + i = \operatorname{trace} \begin{bmatrix} a & b & c \\ d & e & f \\ g & h & i \end{bmatrix}$$

$$\operatorname{curl}(\mathbf{u}) = \det \begin{bmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ ax + by + cz & dx + ey + fz & gx + hy + iz \end{bmatrix}$$

$$= (h - f, c - g, d - b)$$

Div and curl of matrix fields

Consider a vector field

$$\mathbf{u} = \begin{bmatrix} ax + by + cz \\ dx + ey + fz \\ gx + hy + iz \end{bmatrix} = \begin{bmatrix} a & b & c \\ d & e & f \\ g & h & i \end{bmatrix} \begin{bmatrix} x \\ y \\ z \end{bmatrix}$$

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$$\operatorname{div}(\mathbf{u}) = \frac{\partial}{\partial x}(ax + by + cz) + \frac{\partial}{\partial y}(dx + ey + fz) + \frac{\partial}{\partial z}(gx + hy + iz)$$

$$= a + e + i = \operatorname{trace} \begin{bmatrix} a & b & c \\ d & e & f \\ g & h & i \end{bmatrix}$$

$$\operatorname{curl}(\mathbf{u}) = \det \begin{bmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ ax + by + cz & dx + ey + fz & gx + hy + iz \end{bmatrix}$$
$$= (h - f, c - g, d - b)$$

So $\operatorname{curl}(\mathbf{u}) = 0$ if the matrix $\begin{bmatrix} a & b & c \\ d & e & f \\ g & h & i \end{bmatrix}$ is symmetric.

- (a) For a scalar field f in two dimensions, $\text{grad}(f) = \nabla(f) = (f_x, f_y)$
(a vector field).

grad, div and curl in two dimensions

- (a) For a scalar field f in two dimensions, $\text{grad}(f) = \nabla(f) = (f_x, f_y)$
(a vector field).
- (b) For a vector field $\mathbf{u} = (p, q)$ in two dimensions, $\text{div}(\mathbf{u}) = \nabla \cdot \mathbf{u} = p_x + q_y$
(a scalar field).

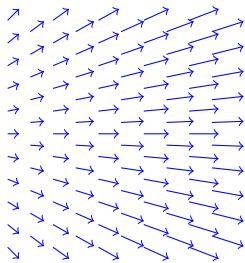
- (a) For a scalar field f in two dimensions, $\text{grad}(f) = \nabla(f) = (f_x, f_y)$
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- (b) For a vector field $\mathbf{u} = (p, q)$ in two dimensions, $\text{div}(\mathbf{u}) = \nabla \cdot \mathbf{u} = p_x + q_y$
(a scalar field).
- (c) For a vector field $\mathbf{u} = (p, q)$ in two dimensions,

$$\text{curl}(\mathbf{u}) = \det \begin{bmatrix} \frac{\partial}{\partial x} & \frac{\partial}{\partial y} \\ p & q \end{bmatrix} = q_x - p_y$$

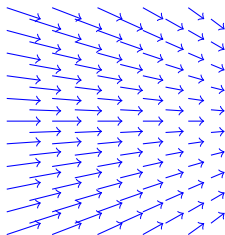
(a *scalar* field, not a vector field as in three dimensions).

Geometric interpretation of $\text{div}(\mathbf{u})$

It works out that the divergence $\text{div}(\mathbf{u}) = \nabla \cdot \mathbf{u}$ is positive when the vectors \mathbf{u} are spreading out, and negative when they are coming together.



diverging: $\nabla \cdot \mathbf{u} > 0$

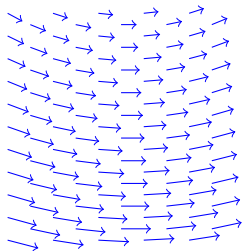


converging: $\nabla \cdot \mathbf{u} < 0$

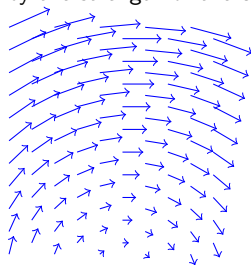
For the velocity field of an incompressible fluid we will have $\nabla \cdot \mathbf{u} = 0$.

Geometric interpretation of $\text{curl}(\mathbf{u})$

In two dimensions, it works out that $\text{curl}(\mathbf{u}) > 0$ in regions where the field is curling anticlockwise, and $\text{curl}(\mathbf{u}) < 0$ in regions where it is curling clockwise, and the absolute value of $\text{curl}(\mathbf{u})$ is determined by the strength of the curling.



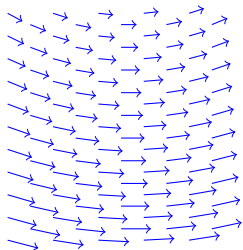
$\text{curl}(\mathbf{u}) > 0$, smaller



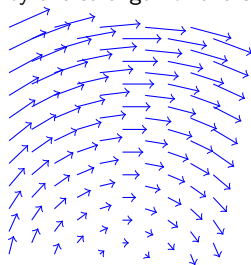
$\text{curl}(\mathbf{u}) < 0$, larger

Geometric interpretation of $\text{curl}(\mathbf{u})$

In two dimensions, it works out that $\text{curl}(\mathbf{u}) > 0$ in regions where the field is curling anticlockwise, and $\text{curl}(\mathbf{u}) < 0$ in regions where it is curling clockwise, and the absolute value of $\text{curl}(\mathbf{u})$ is determined by the strength of the curling.



$\text{curl}(\mathbf{u}) > 0$, smaller



$\text{curl}(\mathbf{u}) < 0$, larger

In three dimensions, the field \mathbf{u} can curl around any axis. In this context, $\text{curl}(\mathbf{u})$ is also a vector field, and it will point along the axis of the curling.

Maxwell's equations

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- ▶ The electric field \mathbf{E} , which is a vector field.

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- ▶ The current density \mathbf{J} , which is also a vector field.
- ▶ The charge density ρ , which is a scalar field.
- ▶ Two constants: $\epsilon_0 \simeq 8.854 \times 10^{-12} F/m^2$ and $\mu_0 \simeq 1.257 \times 10^{-6} Hm^{-1}$.

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The quantities \mathbf{E} , \mathbf{B} , \mathbf{J} and ρ may also depend on time; we write $\dot{\mathbf{E}}$ for $\partial\mathbf{E}/\partial t$ and so on.

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The quantities \mathbf{E} , \mathbf{B} , \mathbf{J} and ρ may also depend on time; we write $\dot{\mathbf{E}}$ for $\partial\mathbf{E}/\partial t$ and so on. The various fields are related by the following equations:

$$\nabla \cdot \mathbf{E} = \rho / \epsilon_0$$

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$$\nabla \times \mathbf{E} = -\dot{\mathbf{B}}$$

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$$\nabla \times \mathbf{B} = \mu_0 \mathbf{J} + \mu_0 \epsilon_0 \dot{\mathbf{E}}.$$

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$$\nabla \cdot \mathbf{E} = \rho / \epsilon_0$$

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$$\nabla \cdot \mathbf{B} = 0$$

$$\nabla \times \mathbf{B} = \mu_0 \mathbf{J} + \mu_0 \epsilon_0 \dot{\mathbf{E}}.$$

This means that:

- ▶ The electric field diverges in regions where there is positive charge, and converges in regions where there is negative charge.

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$$\begin{aligned}\nabla \cdot \mathbf{E} &= \rho / \epsilon_0 & \nabla \times \mathbf{E} &= -\dot{\mathbf{B}} \\ \nabla \cdot \mathbf{B} &= 0 & \nabla \times \mathbf{B} &= \mu_0 \mathbf{J} + \mu_0 \epsilon_0 \dot{\mathbf{E}}.\end{aligned}$$

This means that:

- ▶ The electric field diverges in regions where there is positive charge, and converges in regions where there is negative charge.
- ▶ The magnetic field never diverges or converges.

These involve:

- ▶ The electric field \mathbf{E} , which is a vector field.
- ▶ The magnetic field \mathbf{B} , which is another vector field.
- ▶ The current density \mathbf{J} , which is also a vector field.
- ▶ The charge density ρ , which is a scalar field.
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$$\nabla \cdot \mathbf{E} = \rho / \epsilon_0$$

$$\nabla \times \mathbf{E} = -\dot{\mathbf{B}}$$

$$\nabla \cdot \mathbf{B} = 0$$

$$\nabla \times \mathbf{B} = \mu_0 \mathbf{J} + \mu_0 \epsilon_0 \dot{\mathbf{E}}$$

This means that:

- ▶ The electric field diverges in regions where there is positive charge, and converges in regions where there is negative charge.
- ▶ The magnetic field never diverges or converges.
- ▶ Changing magnetic fields cause the electric field to curl.

These involve:

- ▶ The electric field \mathbf{E} , which is a vector field.
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- ▶ The current density \mathbf{J} , which is also a vector field.
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$$\begin{aligned}\nabla \cdot \mathbf{E} &= \rho / \epsilon_0 & \nabla \times \mathbf{E} &= -\dot{\mathbf{B}} \\ \nabla \cdot \mathbf{B} &= 0 & \nabla \times \mathbf{B} &= \mu_0 \mathbf{J} + \mu_0 \epsilon_0 \dot{\mathbf{E}}.\end{aligned}$$

This means that:

- ▶ The electric field diverges in regions where there is positive charge, and converges in regions where there is negative charge.
- ▶ The magnetic field never diverges or converges.
- ▶ Changing magnetic fields cause the electric field to curl.
- ▶ Currents cause the magnetic field to curl.

These involve:

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- ▶ The current density \mathbf{J} , which is also a vector field.
- ▶ The charge density ρ , which is a scalar field.
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The quantities \mathbf{E} , \mathbf{B} , \mathbf{J} and ρ may also depend on time; we write $\dot{\mathbf{E}}$ for $\partial\mathbf{E}/\partial t$ and so on. The various fields are related by the following equations:

$$\begin{aligned}\nabla \cdot \mathbf{E} &= \rho / \epsilon_0 & \nabla \times \mathbf{E} &= -\dot{\mathbf{B}} \\ \nabla \cdot \mathbf{B} &= 0 & \nabla \times \mathbf{B} &= \mu_0 \mathbf{J} + \mu_0 \epsilon_0 \dot{\mathbf{E}}.\end{aligned}$$

This means that:

- ▶ The electric field diverges in regions where there is positive charge, and converges in regions where there is negative charge.
- ▶ The magnetic field never diverges or converges.
- ▶ Changing magnetic fields cause the electric field to curl.
- ▶ Currents cause the magnetic field to curl. Changing electric fields also cause the magnetic field to curl, but the effect is usually much weaker, because ϵ_0 is small.

Plane wave solution to Maxwell's equations

One class of solutions to Maxwell's equations is as follows.

Plane wave solution to Maxwell's equations

One class of solutions to Maxwell's equations is as follows. Put $c = 1/\sqrt{\mu_0\epsilon_0} \simeq 3 \times 10^8 \text{ms}^{-1}$ (which turns out to be the speed of light)

Plane wave solution to Maxwell's equations

One class of solutions to Maxwell's equations is as follows. Put $c = 1/\sqrt{\mu_0\epsilon_0} \simeq 3 \times 10^8 \text{ms}^{-1}$ (which turns out to be the speed of light), and let α be any constant.

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One class of solutions to Maxwell's equations is as follows. Put $c = 1/\sqrt{\mu_0\epsilon_0} \simeq 3 \times 10^8 \text{ms}^{-1}$ (which turns out to be the speed of light), and let α be any constant. We can take $\mathbf{J} = 0$ and $\rho = 0$ and

$$\mathbf{E} = (0, \sin(\alpha(x - ct)), 0)$$

$$\mathbf{B} = (0, 0, \sin(\alpha(x - ct))/c).$$

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$$\mathbf{E} = (0, \sin(\alpha(x - ct)), 0) \qquad \mathbf{B} = (0, 0, \sin(\alpha(x - ct))/c).$$

We find that

$$\nabla \cdot \mathbf{E} = \frac{\partial}{\partial y} \sin(\alpha(x - ct))$$

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We find that

$$\nabla \cdot \mathbf{E} = \frac{\partial}{\partial y} \sin(\alpha(x - ct)) = 0 = \rho/\epsilon_0 \qquad \dot{\mathbf{E}} = (0, -\alpha c \cos(\alpha(x - ct)), 0)$$

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$$\nabla \cdot \mathbf{B} = \frac{\partial}{\partial z} \sin(\alpha(x - ct))/c = 0 \qquad \dot{\mathbf{B}} = (0, 0, -\alpha \cos(\alpha(x - ct)))$$

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$$\nabla \cdot \mathbf{B} = \frac{\partial}{\partial z} \sin(\alpha(x - ct))/c = 0 \qquad \dot{\mathbf{B}} = (0, 0, -\alpha \cos(\alpha(x - ct)))$$

$$\nabla \times \mathbf{E} = \det \begin{bmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ 0 & \sin(\alpha(x - ct)) & 0 \end{bmatrix}$$

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$$\nabla \cdot \mathbf{B} = \frac{\partial}{\partial z} \sin(\alpha(x - ct))/c = 0 \qquad \dot{\mathbf{B}} = (0, 0, -\alpha \cos(\alpha(x - ct)))$$

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$$\nabla \cdot \mathbf{B} = \frac{\partial}{\partial z} \sin(\alpha(x - ct))/c = 0 \qquad \dot{\mathbf{B}} = (0, 0, -\alpha \cos(\alpha(x - ct)))$$

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Plane wave solution to Maxwell's equations

One class of solutions to Maxwell's equations is as follows. Put $c = 1/\sqrt{\mu_0\epsilon_0} \simeq 3 \times 10^8 \text{ms}^{-1}$ (which turns out to be the speed of light), and let α be any constant. We can take $\mathbf{J} = 0$ and $\rho = 0$ and

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This shows that we do indeed have a solution to the equations.

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This shows that we do indeed have a solution to the equations. It represents an electromagnetic wave of wavelength $1/\alpha$ moving at speed c in the x -direction.

Stationary charged particle

Another solution to Maxwell's equations has $\mathbf{E} = (-xr^{-3}, -yr^{-3}, -zr^{-3})$ with all other fields (\mathbf{B} , \mathbf{J} and ρ) being zero.

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$$\begin{aligned} \nabla \cdot \mathbf{E} &= (-xr^{-3})_x + (-yr^{-3})_y + (-zr^{-3})_z \\ &= 3x^2r^{-5} - r^{-3} + 3y^2r^{-5} - r^{-3} + 3z^2r^{-5} - r^{-3} \\ &= 3(x^2 + y^2 + z^2)r^{-5} - 3r^{-3} = 3r^2r^{-5} - 3r^{-3} = 0. \end{aligned}$$

$$\begin{aligned} \nabla \times \mathbf{E} &= \left((-zr^{-3})_y - (-yr^{-3})_z, (-xr^{-3})_z - (-zr^{-3})_x, (-yr^{-3})_x - (-xr^{-3})_y \right) \\ &= \left(3yzt^{-3} - 3yzt^{-3}, 3xzt^{-3} - 3xzt^{-3}, 3xyr^{-3} - 3xyr^{-3} \right) \end{aligned}$$

Stationary charged particle

Another solution to Maxwell's equations has $\mathbf{E} = (-xr^{-3}, -yr^{-3}, -zr^{-3})$ with all other fields (\mathbf{B} , \mathbf{J} and ρ) being zero. It is clear that $\dot{\mathbf{E}} = 0$ and $\dot{\mathbf{B}} = 0$, so the only equations that we need to check are that $\nabla \cdot \mathbf{E} = 0$ and $\nabla \times \mathbf{E} = 0$. For this we recall that $r_x = x/r$, so $(r^{-3})_x = -3r^{-4}r_x = -3xr^{-5}$. In the same way, we have $(r^{-3})_y = -3yr^{-5}$ and $(r^{-3})_z = -3zr^{-5}$. Using this we find that

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Stationary charged particle

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This shows that we have a solution to the equations, as claimed.

Stationary charged particle

Another solution to Maxwell's equations has $\mathbf{E} = (-xr^{-3}, -yr^{-3}, -zr^{-3})$ with all other fields (\mathbf{B} , \mathbf{J} and ρ) being zero. It is clear that $\dot{\mathbf{E}} = 0$ and $\dot{\mathbf{B}} = 0$, so the only equations that we need to check are that $\nabla \cdot \mathbf{E} = 0$ and $\nabla \times \mathbf{E} = 0$. For this we recall that $r_x = x/r$, so $(r^{-3})_x = -3r^{-4}r_x = -3xr^{-5}$. In the same way, we have $(r^{-3})_y = -3yr^{-5}$ and $(r^{-3})_z = -3zr^{-5}$. Using this we find that

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This shows that we have a solution to the equations, as claimed. This one represents the electric field of a single stationary particle at the origin, with no magnetic field.